# QUANTUM FIELD THEORY I PROBLEMS 2nd semester 2024-2024

Solutions to Srendicki's problems are available online, e.g. at

 $\tt http://https://www.studocu.com/de/document/universitat-heidelberg/theoretische-physik-iv-quantenmechanik/mark-srednicki-quantum-field-theory-problem-solutions/103881937$ 

#### I: Classical field theory

(1) Consider a system of N coupled harmonic oscillators with potential

$$V = \sum_{i=1}^{N} \frac{\kappa}{2} (q_{i+1} - q_i)^2.$$
(1)

Determine the normal coordinates and the eigenvectors of the potential, with boundary conditions  $q_0(t) = q_{N+1}(t) = 0$ . Determine directly the equations of motion and the normal coordinates in the continuum limit without using the Lagrangian, by taking the continuum limit of the result obtained in the discrete case. Solution: e.g.

https://phys.libretexts.org/Bookshelves/Classical.Mechanics/Variational.Principles.in\_Classical.Mechanics\_(Cline)/14%3A\_Coupled\_Linear\_Oscillators/14.10%3A\_Discrete\_Lattice\_Chain

(2) Consider the Lagrangian density

$$\mathcal{L}(x) = \partial_{\mu}\phi^*\partial_{\mu}\phi - m^2|\phi|^2.$$
<sup>(2)</sup>

where  $\phi$  is a complex classical field. Derive the classical equations of motion and solve them using the method of normal coordinates.

(3) Optional: problems 5.2a-d, 5.4 and 5.5 of Radovanović

II: The Poincaré group, symmetries and Noether's theorem

(4) Show that requiring invariance of the metric upon Lorentz transformations

$$\Lambda^{\mu}{}_{\nu}\Lambda^{\rho}{}_{\sigma}g^{\nu\sigma} = g^{\mu\rho},\tag{3}$$

and writing the infinitesimal transformation in terms of its generators  $J_{\rho\sigma}$ 

$$\Lambda^{\mu}{}_{\nu} = g^{\mu}{}_{\nu} + \frac{i}{2}\omega^{\rho\sigma}J_{\rho\sigma}{}^{\mu}{}_{\nu}, \qquad (4)$$

completely fixes the explicit form of the generators, viewed as  $4 \times 4$  Lorenz matrices  $(J_{\rho\sigma})^{\mu}{}_{\nu}$ , and determine their explicit expression.

Solution: see Eq.s (13-19) of the notes on the website.

(5) Derive the Lorentz algebra, i.e. the commutation relations

$$[J^{\mu\nu}, J^{\rho\sigma}] = -i\left(g^{\mu\sigma}J^{\nu\rho} + g^{\nu\rho}J^{\mu\sigma} - g^{\mu\rho}J^{\nu\sigma} - g^{\nu\sigma}J^{\mu\rho}\right)$$
(5)

- (a) using the explicit form of the generators which was found in problem (3);
- (b) from the relation  $D(\Lambda')^{-1}D(\Lambda)D(\Lambda') = D(\Lambda'^{-1}\Lambda\Lambda')$ , which holds for any Lorentz representation  $D(\Lambda)$  in the case of infinitesimal transformations (see problem 1.11-b of Radovanović).
- (6) Optional: problems 1.16 and 1.17 of Radovanović

(7) Show that the Lagangian density

$$\mathcal{L} = \frac{1}{2} \left( \partial_{\mu} \phi_1 \partial_{\mu} \phi_1 + \partial_{\mu} \phi_2 \partial_{\mu} \phi_2 \right) - \frac{m^2}{2} \left( \phi_1^2 + \phi_2^2 \right) - \lambda \left( \phi_1^2 + \phi_2^2 \right)^2 \tag{6}$$

is invariant under the transformation

$$\begin{pmatrix} \phi_1' \\ \phi_2' \end{pmatrix} = R(\theta) \begin{pmatrix} \phi_1 \\ \phi_2 \end{pmatrix}$$
(7)

where  $R(\theta)$  is a rotation matrix by an angle  $\theta$ . Determine the Noether current and charge. Solution: see problem 5.11 of Radovanović.

- (8) Optional: problems 5.12, 5.14b, 5.15, and 5.17 of Radovanović
- (9) Consider the Maxwell Langrangian

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu},\tag{8}$$

with

$$F_{\mu\nu} = \partial_{\mu}A_{\nu} - \partial_{\nu}A_{\mu}.$$
(9)

Determine the energy-momentum tensor  $T'^{\mu\nu}$  explicitly in terms of the fields  $a^{\mu}$  through Noether's theorem, and show that it is not symmetric. Determine explicitly the prepotential  $A^{\rho\mu\nu} = -A^{\mu\rho\nu}$  such that

$$T^{\mu\nu} = T^{\prime\mu\nu} + \partial_{\rho} A^{\rho\mu\nu} \tag{10}$$

is symmetric and given by the so-called Belinfante form

$$T^{\mu\nu} = -\frac{1}{2} \left( F^{\mu\alpha} F^{\nu}{}_{\alpha} - \mathcal{L}g^{\mu\nu} \right).$$
 (11)

Solution: see problem 5.18 of Radovanović.

(10) Optional: problem 5.19 of Radovanović.

# III: Canonical field quantization of scalar fields

(11) Derive the canonical commutation relation for the scalar field operator

$$[\phi(\vec{x}), \pi(\vec{x}')] = i\delta^{(3)}(\vec{x} - \vec{x}') \tag{12}$$

from the expression of the field operators  $\phi$  and  $\phi$  in terms of creation and annihilation operators, and the commutation relation satisfied by the latter.

(12) Show that the momentum operator for a real scalar quantum field, namely

$$\vec{P} = -\int d^3x \dot{\phi}(\vec{x}, t) \vec{\nabla} \phi(\vec{x}, t)$$
(13)

generates translations of the field operator, i.e.

$$[\vec{P}, \phi(\vec{x}, t)] = -i\vec{\nabla}\phi(\vec{x}, t).$$
(14)

(13) Optional: problems 7.6 b-f, 7.7 and 7.8 of Radovanović.

## IV: Vector and spinor fields

(14) Show that the spin operator for the electromagnetic field has the form

$$s^{i} = \epsilon^{ijk} \int d^{3}x E_{j} A_{k} \tag{15}$$

$$= \int \frac{d^3k}{(2\pi)^3} \sum_{s=\pm} \left( \epsilon^{ijk} \varepsilon^{s*}_{\vec{k}j} \varepsilon^s_{\vec{k}k} \right) a^{s\dagger}_{\vec{k}} a^s_{\vec{k}}.$$
(16)

(15) Prove that if  $\gamma^{\mu}$  are  $n \times n$  matrices satisfying the Clifford algebra

$$\{\gamma^{\mu}, \gamma^{\nu}\} = 2g^{\mu\mu}I,\tag{17}$$

where I is the n-dimensional identity matrix, then the matrices

$$\sigma^{\mu\nu} = \frac{i}{4} [\gamma^{\mu}, \gamma^{\nu}] \tag{18}$$

satisfy the Lorentz algebra, as given in problem 4, Eq. (4).

(16) Prove that the matrices  $\sigma^{\mu\nu}$  Eq. (16) satisfy the equation

$$\Lambda_{1/2}^{-1}\gamma^{\mu}\Lambda_{1/2} = \Lambda^{\mu}{}_{\nu}\gamma^{\nu}, \qquad (19)$$

where

$$\Lambda_{1/2} = \exp{-\frac{i}{2}\omega_{\alpha\beta}\sigma^{\alpha\beta}} \tag{20}$$

$$\Lambda^{\mu}{}_{\nu} = \left(\exp -\frac{i}{2}\omega_{\alpha\beta}(J^{\alpha\beta})^{\mu}{}_{\nu}$$
(21)

and  $J^{\alpha\beta}$  are the generators of the Lorentz group (see problem 4)

$$\left(J^{\alpha\beta}\right)^{\mu}{}_{\nu} = i \left(g^{\mu\alpha}g^{\beta}{}_{\nu} - g^{\mu\beta}g^{\alpha}{}_{\nu}\right).$$
<sup>(22)</sup>

It is sufficient to prove the result in the case of infinitesimal transformations.

- (17) Write the Dirac equations satisfied by the left and right component of a generic Dirac field  $\psi$ , defined as  $\psi_L \equiv \frac{\mathbb{I} \gamma + 5}{2}\psi$ ;  $\psi_R \equiv \frac{\mathbb{I} + \gamma + 5}{2}\psi$ . Determine explicitly the form of the  $\gamma_5$  matrix in the Weyl representation of Dirac matrices.
- (18) Show that if the solutions of the Dirac equations are normalized as

$$\bar{u}_r(p)u_s(p) = 2m\delta_{rs} \tag{23}$$

then

$$u_r(p)^{\dagger} u_s(p) = 2E\delta_{rs}, \tag{24}$$

where  $E = p^0$ .

- (19) Optional: problems 4.6, 4.10, 4.11 and 5.13 of Radovanović.
- (20) Define creation and annihilation operators  $b_s(p)$ ,  $b_s^{\dagger}(p)$ ,  $d_s(p)$  and  $d_s^{\dagger}(p)$  for the Dirac field according to

$$\psi(x) = \int \frac{d^3p}{(2\pi)^3} \frac{1}{\sqrt{2E}} \sum_s \left( u^s(p) b_s(p) e^{-ipx} + v^s(p) d_s^{\dagger}(p) e^{ipx} \right)$$
(25)

where u and v are solution to the momentum-space Dirac equation, and determine in term of these operators the expression of the Dirac Hamiltonian

$$H = \int d^3x \bar{\psi} \left( i \vec{\gamma} \cdot \vec{\nabla} + m \right) \psi(x).$$
<sup>(26)</sup>

Solution: see problem 8.2b of Radovanović.

(21) Optional: problems 5.2e, 8.2a, c 8.3 and 8.9 Radovanović.

### V: Path integral and Green functions

(22) Determine the classical action for a one-dimensional free particle with fixed initial and final conditions. Determine explicitly the matrix element of the time-evolution operator

$$K(q't';q,t) = \langle q',t'|e^{-iH(t',-t)}|q,t\rangle$$
(27)

and show that is equal to the exponential of the classical action, times a factor of i. Solution: see problem 6.1b,c of Srednicki.

(23) Show that the ground-state wave functional for a free real scalar field such that

$$\hat{\phi}(x)|\phi(x)\rangle = \phi(x)|\phi(x)\rangle$$
(28)

has the form

$$\langle \phi(x)|0\rangle = \mathcal{N}\exp{-\frac{1}{2}\int \frac{d^3k}{(2\pi)^3}\phi^*(k)\sqrt{\vec{k}|^2 + m^2}\phi(k)},$$
 (29)

where  $\mathcal{N}$  is a suitable normalization. Solution: see problem 8.8 of Srednicki.

(24) Compute, using the generating functional, the four-point function

$$G^{(4)}(x_1, x_2, x_3, x_4) = \langle 0 | T \left[ \phi(x_1) \phi(x_2) \phi(x_3) \phi(x_4) \right] | 0 \rangle$$
(30)

for a free real scalar field.

- (25) Optional: problems 6.4, 6.5, 6.9, 6.10 of Radovanović.
- (26) Prove that if  $\theta_i$  are (complex) components of an *n*-dimensional vector of Grassmann numbers and B is an  $n \times n$  matrix then

$$\int \prod_{i} d\theta_{i}^{*} d\theta_{j} e^{-\theta_{i}^{*} B_{ij} \theta_{j}} = B_{ij}^{-1} \det B.$$
(31)

(27) Compute explicitly the two-point function for a Dirac field

$$G^{(2)}(x_1, x_2) = \langle 0 | T\psi(x_1)\bar{\psi}(x_2) \rangle$$
(32)

by expressing the field in terms of creation and annihilation operators and using the anticommutation relations.

Solution: see problem 8.1 of Radovanović.

# VI: The reduction formula and Feynman rules

(28) Prove that if one defines

$$a_{\vec{k}}(t) \equiv i \int \frac{d^3x}{\sqrt{2\omega}} e^{i(\omega t - \vec{k} \cdot \vec{x})} \frac{\vec{d}}{dt} \phi(\vec{x}, t)$$
(33)

then

$$\phi(x) = \int \frac{d^3k}{(2\pi)^3 \sqrt{2\omega}} \left( e^{-ikx} a_{\vec{k}}(t) + e^{ikx} a_{\vec{k}}^{\dagger}(t) \right).$$
(34)

(29) Prove the reduction formula for a fermion field, namely, for incoming and outgoing particles,

$$\langle p_1 \dots p_n | k_1 \dots k_m \rangle = \int d^4 x_1 \dots d^4 x_n d^4 y_1 \dots d^4 y_m e^{i(p_1 x_1 + \dots + p_n x_n - k_1 y_1 - \dots - k_m y_m)}$$
(35)

$$\bar{u}(p_1)S^{-1}(x_1)\dots\bar{u}(p_n)S^{-1}(x_n)\langle T\left(\bar{\psi}(x_1)\dots\bar{\psi}(x_n)\psi(y_1)\dots\psi(y_m)\right)\rangle$$
(36)

$$S^{-1}(y_1)u(k_1)\dots S^{-1}(y_m)u(k_m), (37)$$

and  $u \leftrightarrow v$ ,  $\bar{u} \leftrightarrow \bar{v}$ ,  $k \leftrightarrow -k$  for outgoing or incoming antiparticles respectively (an outgoing particle is like an incoming antiparticle and conversely).